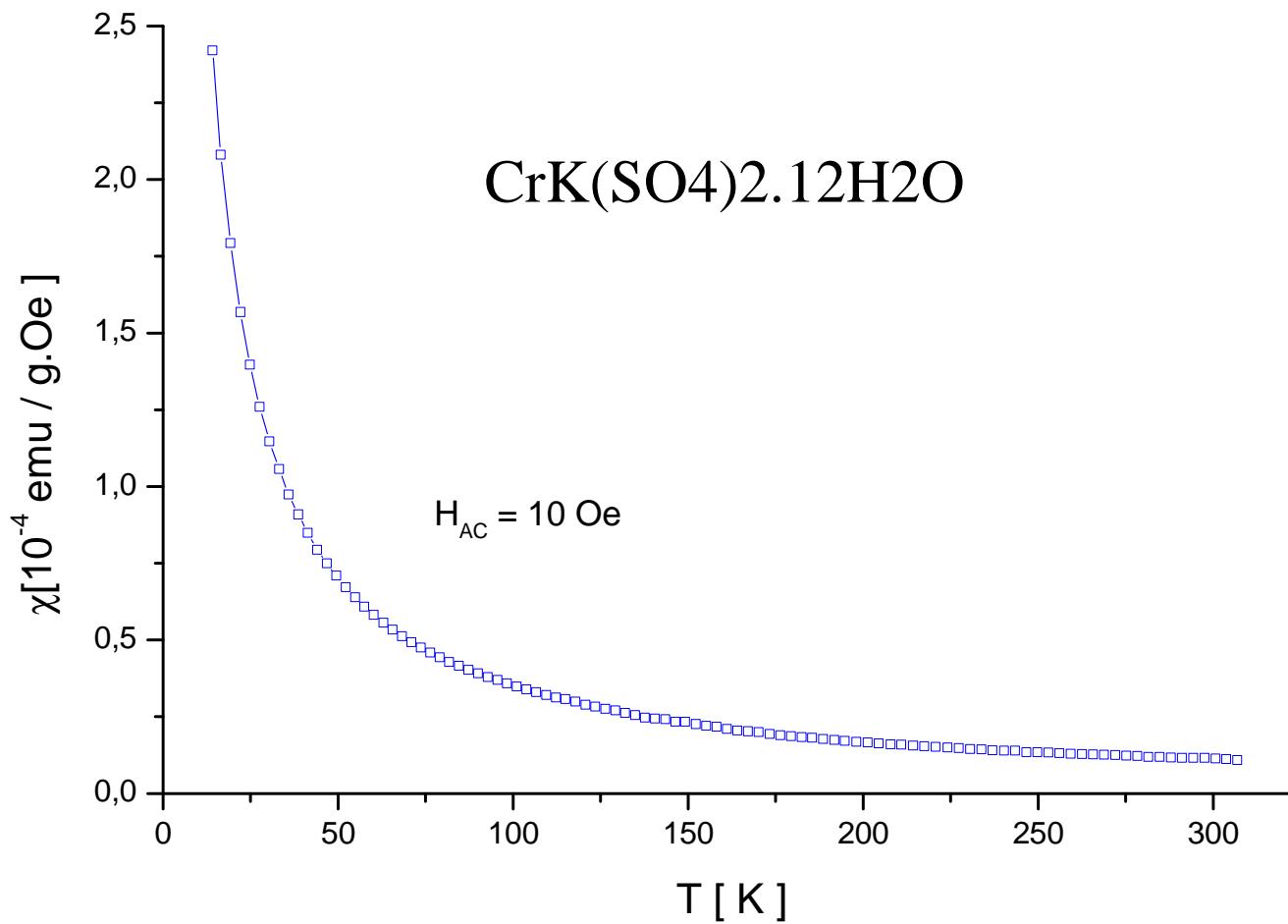


Ejercicio 6



CrK(SO4)2.12H2O

$$\chi_p = M / H \approx \frac{N_{Cr} \mu_0 \mu_B^2 g^2 s(s+1)}{3kT} = \frac{N_{Cr} \mu_0 \mu_B^2 g^2}{3k} s(s+1) \frac{1}{T} = \frac{C}{T}$$

$$C = \frac{N_{Cr} \mu_0 \mu_B^2 g^2}{3k} s(s+1) = 0,0202$$

$$\frac{N_{Cr} \mu_0 \mu_B^2 g^2}{3k} = \frac{N_{Cr} \cdot 4\pi \times 10^{-7} \cdot (9.27 \times 10^{-24})^2 \cdot 4}{3 \cdot 1.38 \times 10^{-23}}$$

$$1m_{Cr} = 51.9961 \text{g/mol}$$

$$1m_K = 39.0983 \text{g/mol}$$

$$2m_S = 2 \times 32.065 \text{g/mol}$$

$$20m_O = 20 \times 15.9994 \text{g/mol}$$

$$24m_H = 24 \times 1.00794 \text{g/mol}$$

$$m_{\text{CrK(SO}_4\text{)}_2\text{.12H}_2\text{O}} = 499,40296 \text{g/mol}$$

$$v_{\text{CrK(SO}_4\text{)}_2\text{.12H}_2\text{O}} = 272.898 \text{cm}^3/\text{mol}$$

$$v_{\text{CrK(SO}_4\text{)}_2\text{.12H}_2\text{O}} = 2,729 \times 10^{-4} \text{m}^3/\text{mol}$$

$$N_{Cr} = 6.02 \times 10^{23} \text{át}/2,729 \times 10^{-4} \text{m}^3 = 2.206 \times 10^{27} \text{át/m}^3$$

$$\frac{N_{Cr}\mu_0\mu_B^2g^2}{3k} = \frac{2.206 \times 10^{27} \cdot 4\pi \times 10^{-7} \cdot (9.27 \times 10^{-24})^2 \cdot 4}{3 \cdot 1.38 \times 10^{-23}} = 0.023$$

$$s(s+1) = \frac{0,0202}{0,023} = 0,876$$

$$s^2 + s - 0,876 = 0$$

$$s = \frac{-1 \pm \sqrt{1+3.506}}{2} = \begin{cases} 0.561 \\ -1.561 \end{cases}$$

$$[Cr^0] = \dots 3d^4 4s^2$$

$$[Cr^{+2}] = \dots 3d^4$$

$$[Cr^{+3}] = \dots 3d^3 \Rightarrow s = 3/2$$

$$[Cr^{+4}] = \dots 3d^2$$

$$[Cr^{+5}] = \dots 3d^1 \Rightarrow s = 1/2$$

$$B_{1/2}(x) = \tanh(x)$$

$$B_s(x)\!=\!\frac{2s+1}{2s}\coth\!\left(\frac{2s+1}{2s}x\right)\!-\!\frac{1}{2s}\coth\!\left(\frac{x}{2s}\right)$$

$$B_{1/2}(x)\!=\!2\coth(2x)\!-\!\coth(x)$$

$$B_{1/2}(x)\!=\!2\frac{e^{2x}+e^{-2x}}{e^{2x}-e^{-2x}}\!-\!\frac{e^x+e^{-x}}{e^x-e^{-x}}$$

$$B_{1/2}(x)\!=\!\frac{1}{e^x-e^{-x}}\!\left[2\frac{e^{2x}+e^{-2x}}{e^x+e^{-x}}\!-\!\left(e^x+e^{-x}\right)\right]$$

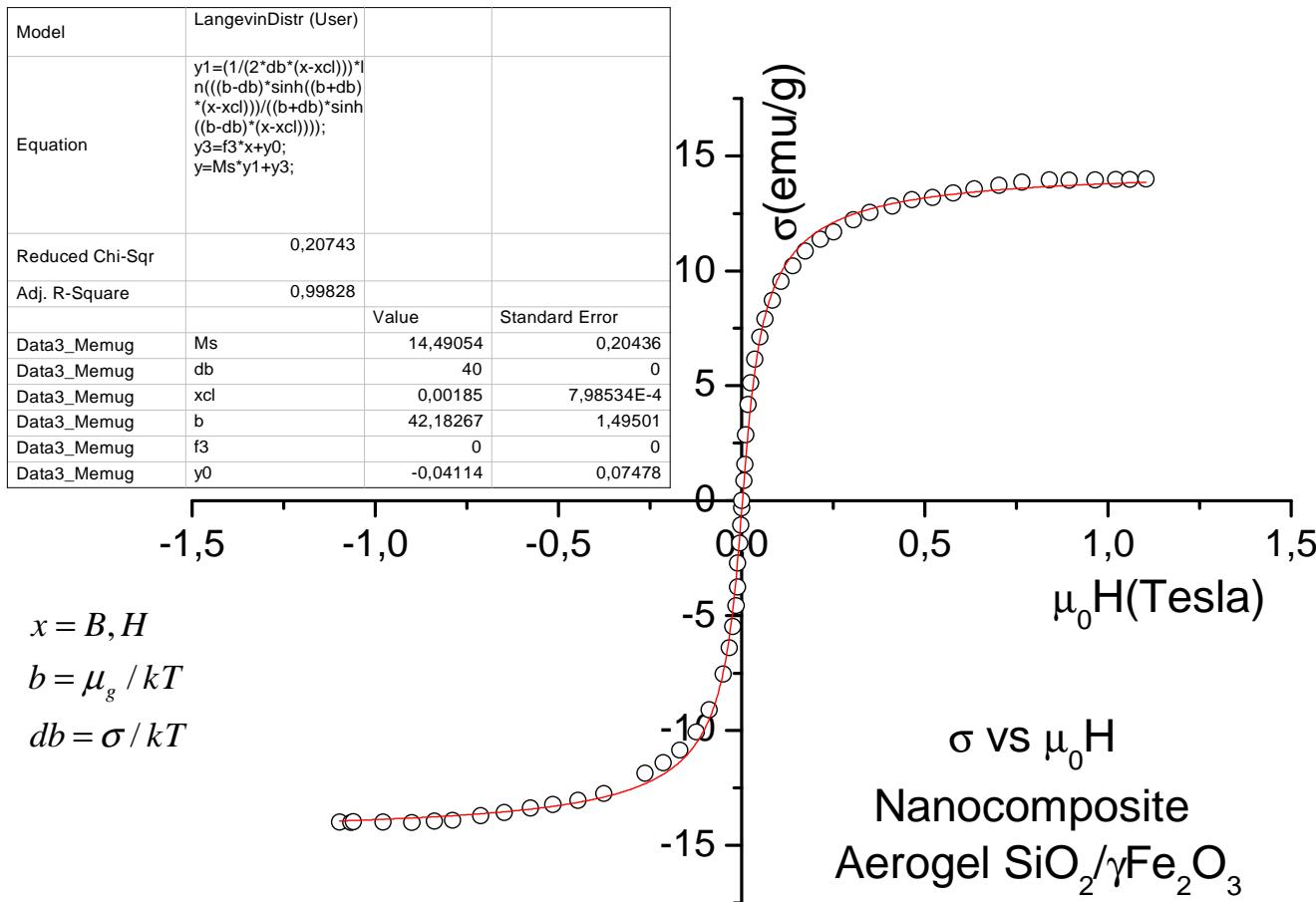
$$B_{1/2}(x)\!=\!\frac{1}{e^x-e^{-x}}\!\left[\frac{2e^{2x}+2e^{-2x}-\left(e^x+e^{-x}\right)^2}{e^x+e^{-x}}\right]$$

$$B_{1/2}(x)\!=\!\frac{1}{e^x-e^{-x}}\!\left[\frac{e^{2x}-2+e^{-2x}}{e^x+e^{-x}}\right]$$

$$B_{1/2}(x)\!=\!\frac{1}{e^x-e^{-x}}\frac{\left(e^x-e^{-x}\right)^2}{e^x+e^{-x}}$$

$$B_{1/2}(x)\!=\!\frac{e^x-e^{-x^2}}{e^x+e^{-x}}=\tanh(x)$$

Ejercicio 9



Análisis de M vs H con Langevin distribuida

Para una muestra monodispersa

$$M = N\mu L(\alpha\mu) \quad N = \frac{\text{part.}}{\text{masa}}; \quad L \equiv \text{Langevin}; \quad \alpha = \frac{B}{kT}$$

Para una muestra polidispersa

$$M = \frac{N \int_0^\infty \mu L(\alpha\mu) P(\mu) d\mu}{\int_0^\infty P(\mu) d\mu}; \quad P(\mu) \equiv \text{distribución de momentos}$$

Momento medio

$$\langle \mu \rangle = \frac{\int_0^\infty \mu P(\mu) d\mu}{\int_0^\infty P(\mu) d\mu} \Rightarrow \int_0^\infty P(\mu) d\mu = \frac{1}{\langle \mu \rangle} \int_0^\infty \mu P(\mu) d\mu$$

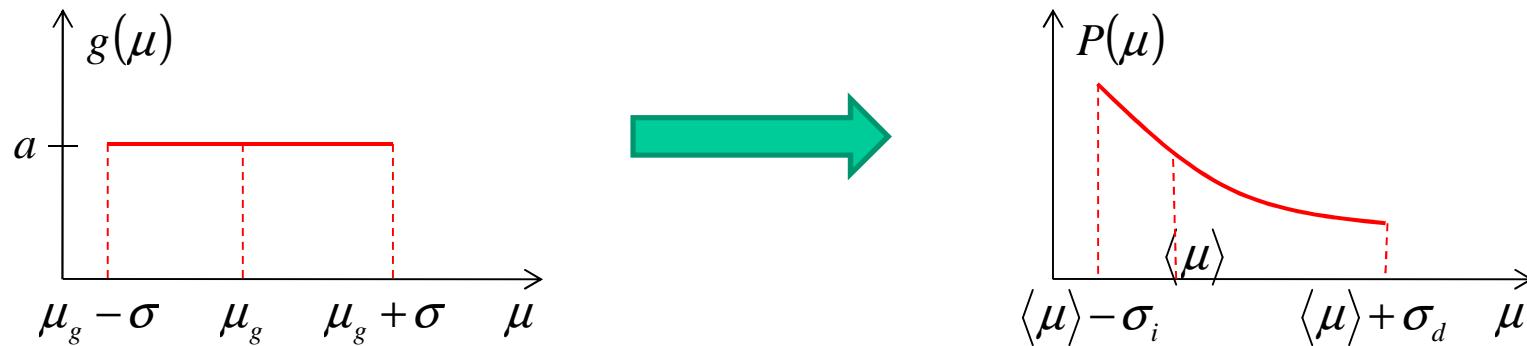


$$M = \frac{N \langle \mu \rangle \int_0^\infty \mu L(\alpha\mu) P(\mu) d\mu}{\int_0^\infty \mu P(\mu) d\mu}$$

Llamando $g(\mu) = \mu P(\mu)$

$$M = \frac{N\langle\mu\rangle \int_0^\infty L(\alpha\mu)g(\mu)d\mu}{\int_0^\infty g(\mu)d\mu} \quad (1)$$

Para poder resolver analíticamente:



$$\begin{aligned} \int_0^\infty P(\mu)d\mu &= a \int_{\mu_1}^{\mu_2} \frac{d\mu}{\mu} = a \ln\left(\frac{\mu_2}{\mu_1}\right) \\ \int_0^\infty \mu P(\mu)d\mu &= \int_0^\infty g(\mu)d\mu = a \int_{\mu_1}^{\mu_2} d\mu = 2a\sigma \end{aligned} \quad \rightarrow \quad \langle\mu\rangle = \frac{\int_0^\infty \mu P(\mu)d\mu}{\int_0^\infty P(\mu)d\mu} = \frac{2\sigma}{\ln\left(\frac{\mu_g + \sigma}{\mu_g - \sigma}\right)}$$

La resolución de (1) es

$$M = \frac{N\langle\mu\rangle}{2\sigma\alpha} \ln \left[\frac{(\mu_g - \sigma) \sinh\{(\mu_g + \sigma)\alpha\}}{(\mu_g + \sigma) \sinh\{(\mu_g - \sigma)\alpha\}} \right] \quad N\langle\mu\rangle = M_s$$

A la expresión anterior le agregamos una componente ferro para tener en cuenta la coercitividad. En un próximo intento trataremos con dos Langevin. Se incluyó y_0 a fin de tener en cuenta un eventual desplazamiento vertical.

La coercitividad xcl se fijó en cero.

$$\gamma = Ms * (((1-f2)/(2*db*(x-xcl))) * \ln(((b-db)*\sinh((b+db)*(x-xcl)))/((b+db)*\sinh((b-db)*(x-xcl)))) + f2*(2/\pi)*\text{atan}((x/xcs-1)*\tan(\pi*S/2)) + f3*x + y0$$

$$x = B, H$$

$$xcs = H_c^{FM}$$

$$b = \mu_g / kT$$

$$S = M_R / M_S^{FM}$$

$$db = \sigma / kT$$

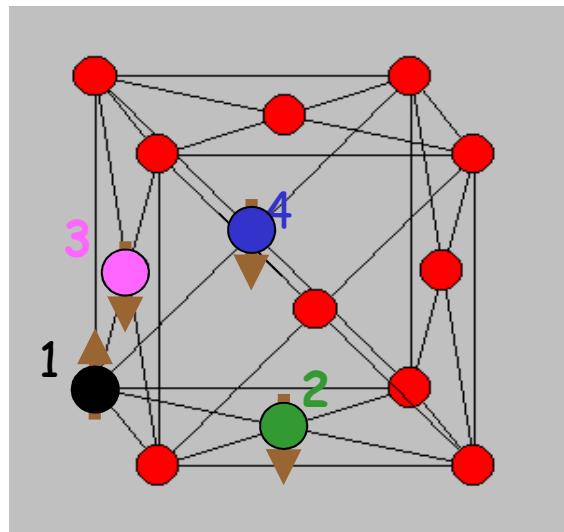
$$f_3 = \text{masa} \cdot \chi_{PM}$$

$$M_s = N_m^{SP} \langle \mu \rangle + M_s^{FM}$$

Intercambio (c)

Ferro y Antiferromagnetismo. Teoría del campo molecular

Red fcc



AF: Primeros vecinos con momentos antiparalelos a "1"

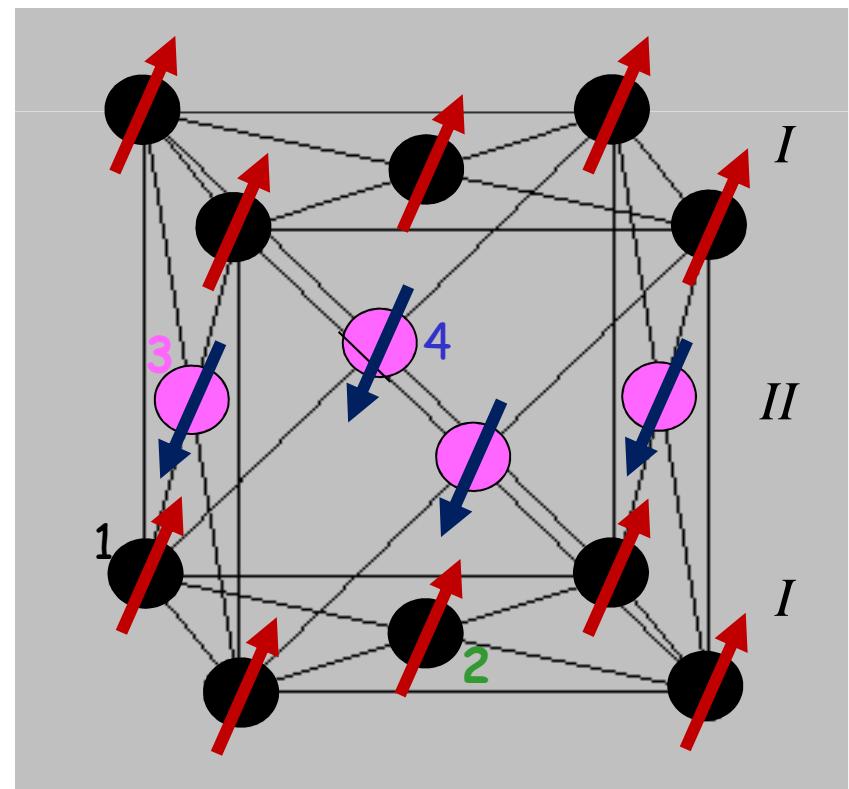
↓
Frustración:
"2" y "4" no
son
antiparalelos

Cómo establecer el orden antiferro

?

Se subdivide la red en 4 subredes, 1, 2, 3, 4

- 1, puntos (k, l, m)
 - 2, puntos $(k+1/2, l+1/2, m)$
 - 3, puntos $(k+1/2, l, m +1/2)$
 - 4, puntos $(k, l +1/2, m +1/2)$
- 2 subredes I, II



Ferro y Antiferromagnetismo. Teoría del campo molecular

Red fcc

$$C = \frac{4Ns(s+1)\mu_0g^2\mu_B^2}{3k}$$

$$\chi_{ini} = \frac{C}{(T + \Theta)} \quad \Theta_{afm} = -\frac{8s(s+1)J}{k}$$

$$\Theta_{fm} = -\frac{8s(s+1)J}{3k}$$

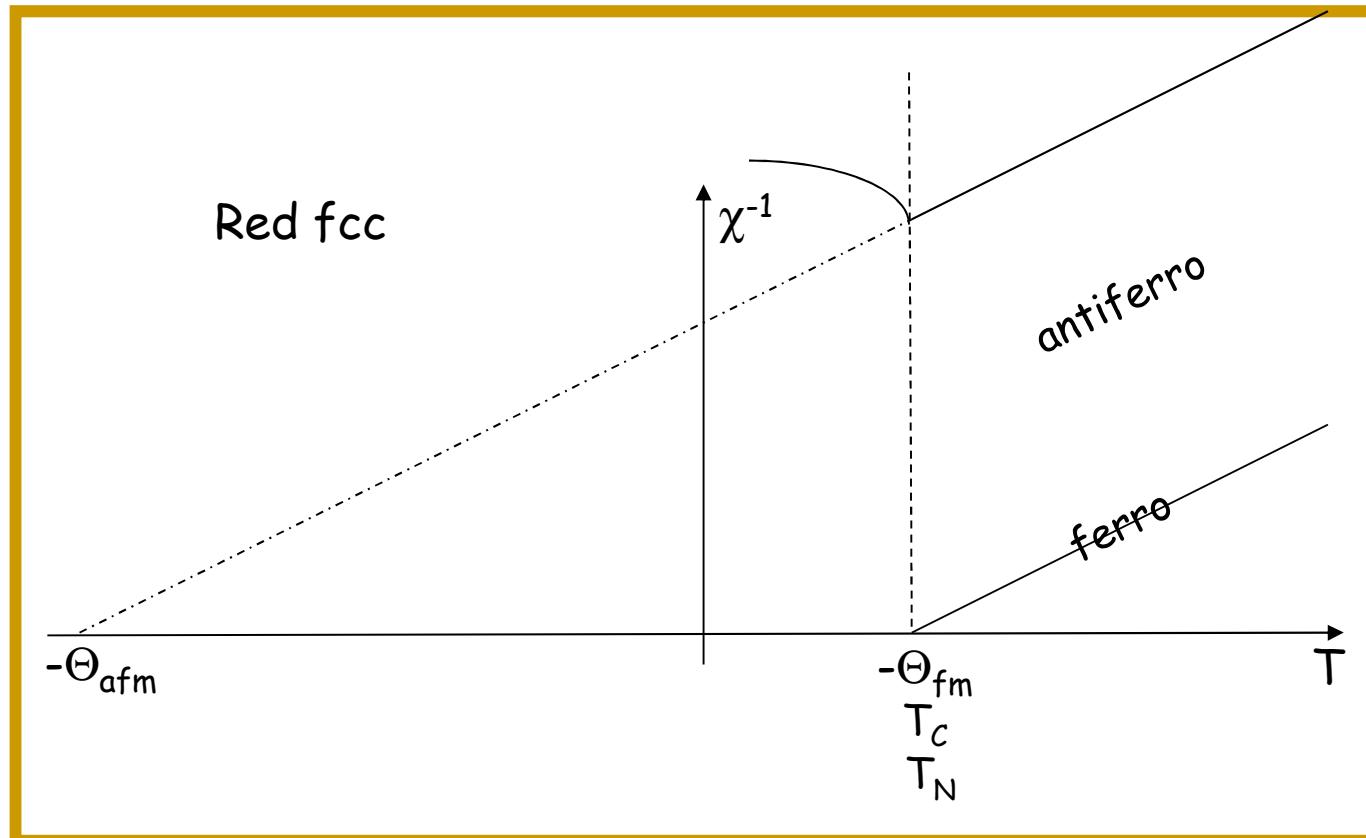
$$J > 0 \quad \rightarrow \quad T_C = \frac{8s(s+1)J}{3k} > 0$$

$$J < 0 \quad \rightarrow \quad T_N = -\frac{8s(s+1)J}{3k} > 0$$

Ferro y Antiferromagnetismo. Teoría del campo molecular

Red fcc

Resumen. Temperaturas características



$$\chi_{ini} = \frac{C}{(T + \Theta)}$$

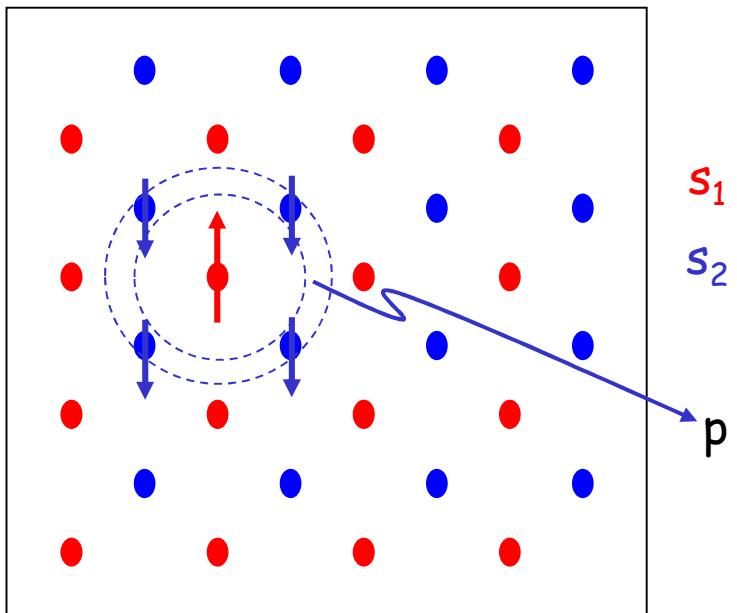
$$\Theta_{afm} = -\frac{8s(s+1)J}{k}$$

$$\Theta_{fm} = -\frac{8s(s+1)J}{3k}$$

$$T_c = \frac{8s(s+1)J}{3k}$$

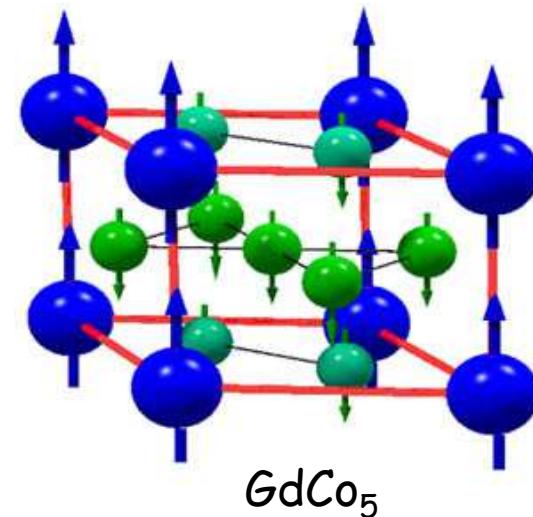
$$T_N = -\frac{8s(s+1)J}{3k}$$

Ferrimagnetismo. Teoría del campo molecular



$$s_1 \neq s_2$$

Ferrimagneto: momento
diferente en cada subred



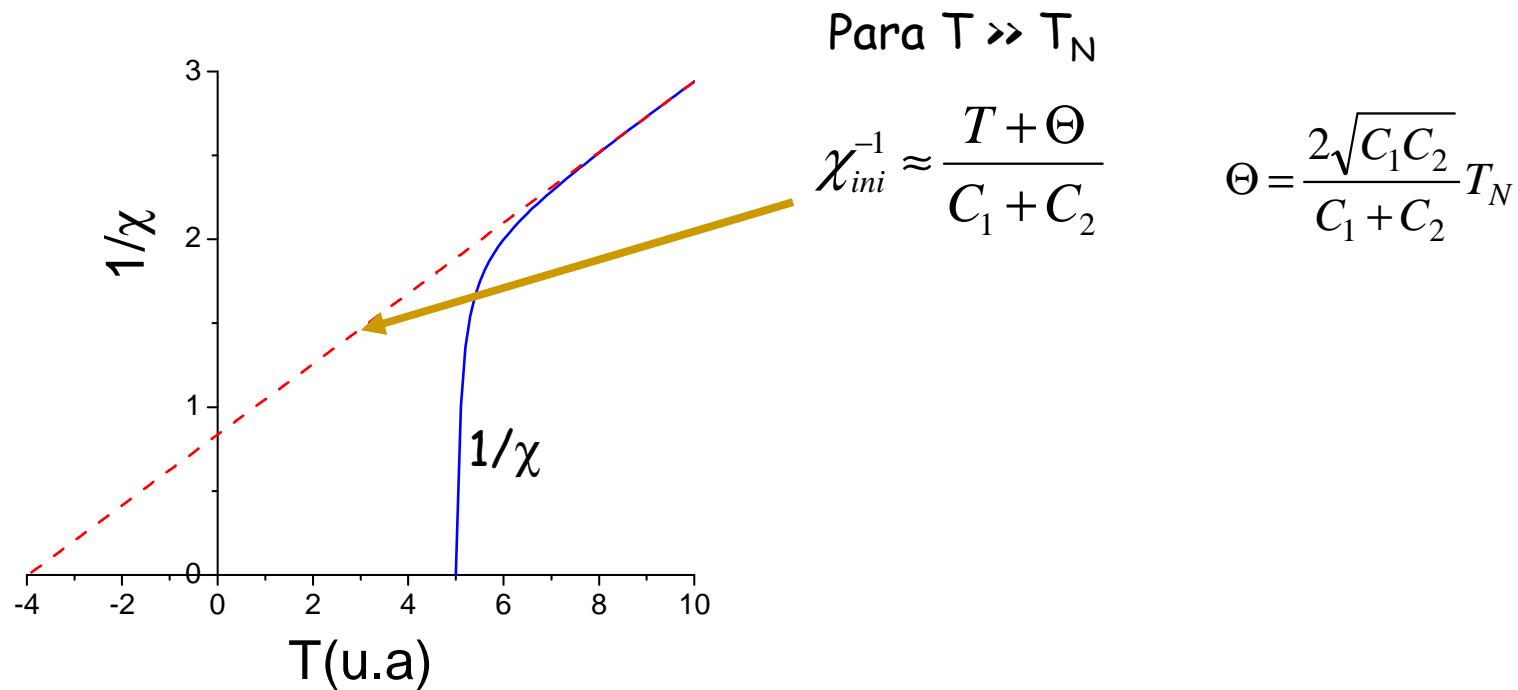
$GdCo_5$
subredes diferentes

Ferrimagnetismo. Teoría del campo molecular

Procedimientos similares a los usados en los casos ferro y antiferro:

$$\chi_{ini} = \frac{(C_1 + C_2)T - 2\sqrt{C_1 C_2 \Theta_1 \Theta_2}}{T^2 - \Theta_1 \Theta_2} = \frac{(C_1 + C_2)T - 2\sqrt{C_1 C_2} T_N}{T^2 - T_N^2}$$

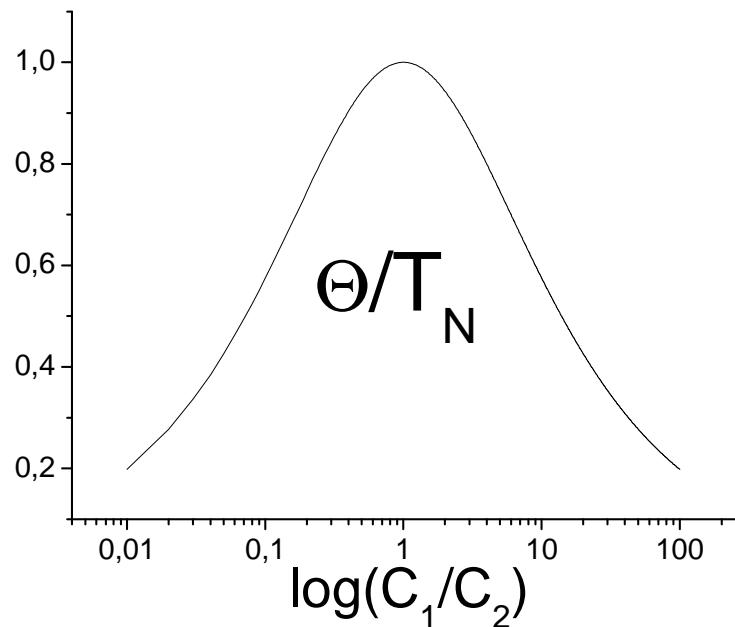
$$C_i = \frac{Ns_i(s_i+1)\mu_0 g_i^2 \mu_B^2}{3k} \quad \Theta_i = \frac{2ps_i(s_i+1)J}{3k} \quad T_N = \sqrt{\Theta_1 \Theta_2}$$



Ferrimagnetismo. Teoría del campo molecular

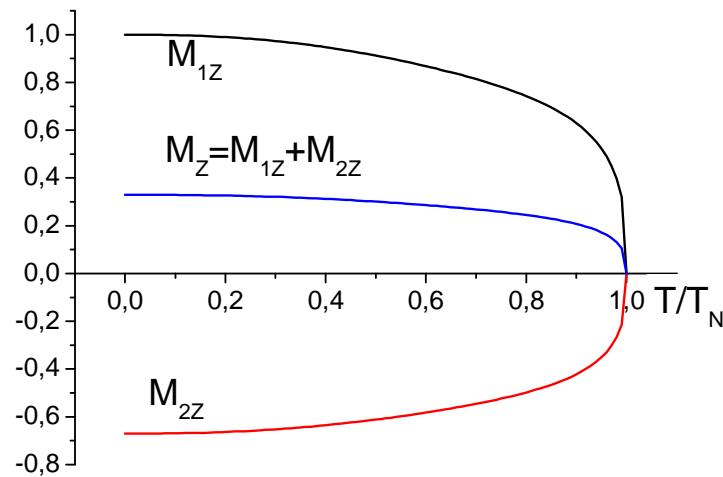
La relación entre Θ y T_N depende de la relación entre C_1 y C_2 :

$$\frac{\Theta}{T_N} = \frac{2\sqrt{C_1 C_2}}{C_1 + C_2} \quad \longrightarrow \quad \begin{aligned} \Theta &\leq T_N \\ \Theta &= T_N \text{ sólo si } C_1 = C_2 \Rightarrow \text{antiferro} \end{aligned}$$

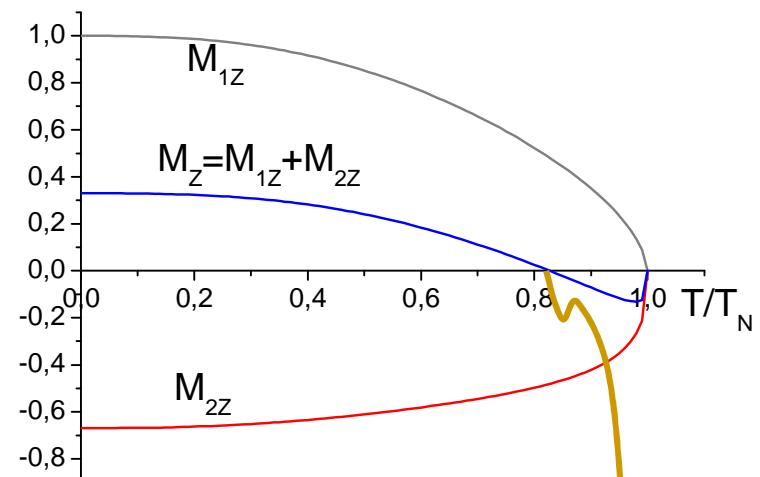


Magnetización de un ferrimagneto

Sub-redes con comportamiento térmico similar

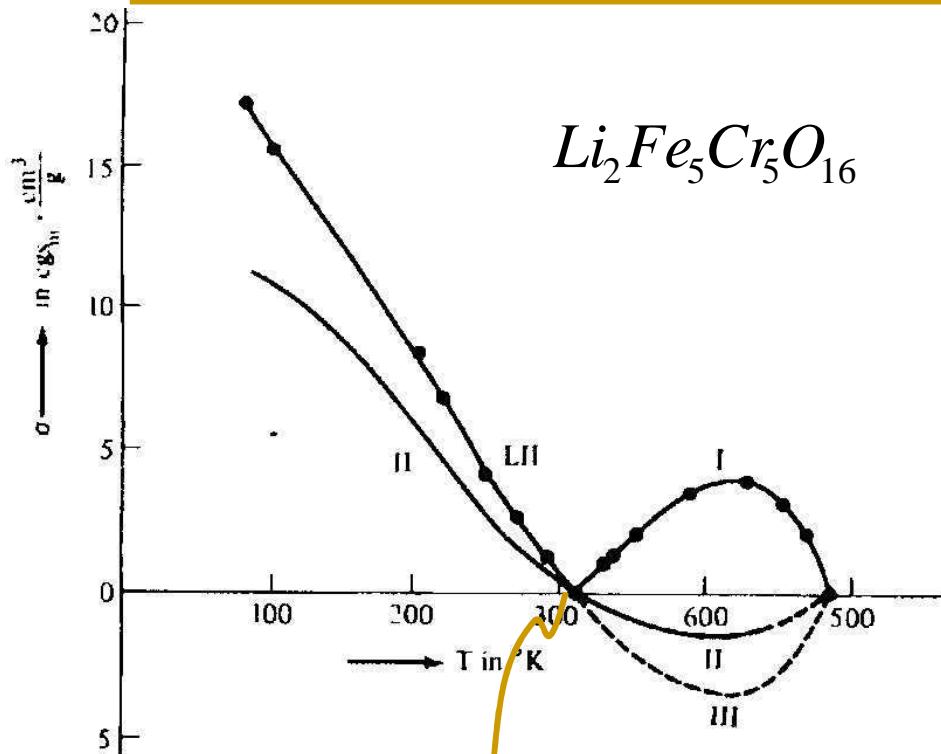


Sub-redes con diferente comportamiento térmico



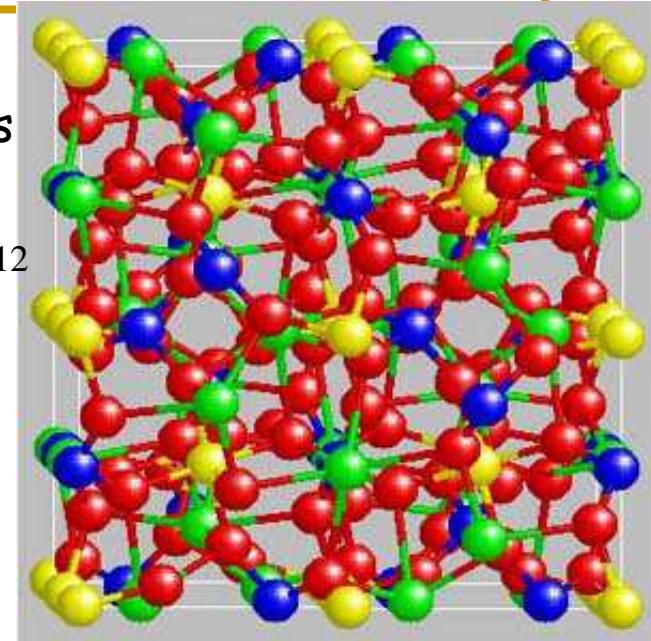
Punto de compensación

Ferrimagnetos reales

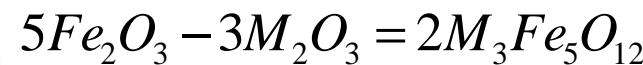


Punto de compensación

garnets



garnets de tierras raras



$\text{Gd}, \text{Tb}, \text{Dy}, \text{Ho}$
 $\text{Er}, \text{Tm}, \text{Yb}, \text{Lu}, \text{Y}$

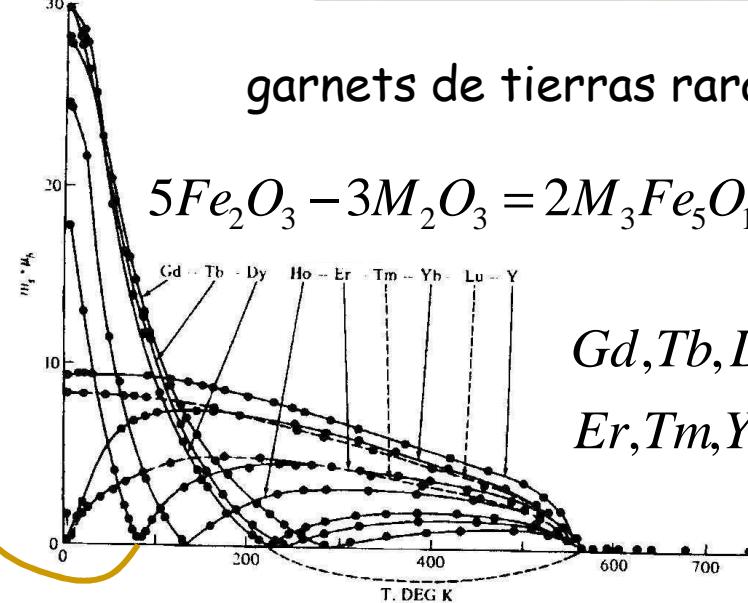
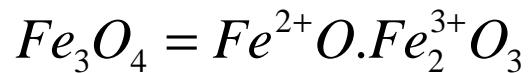
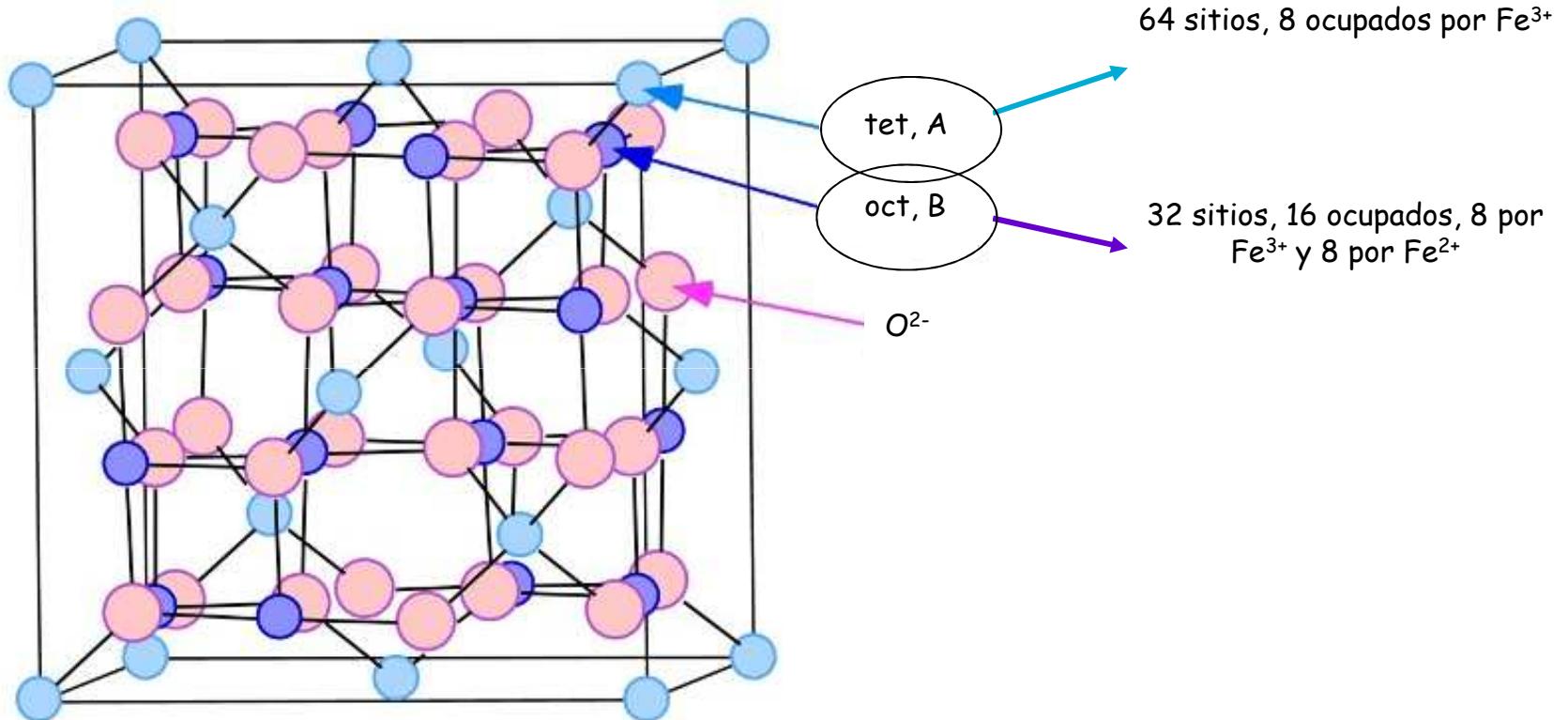


Figure 11.13-Temperature dependencies of the magnetic moments of the rare earth garnets. Note the occurrence of a compensation point in several of the curves. From Bertaut F., and Pauthenet,R., Proc. IEE, 104, Suppl.#5, 261,(1957)

Ferrimagnetos reales

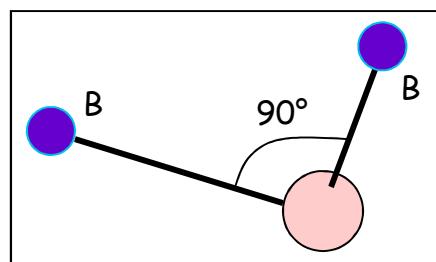
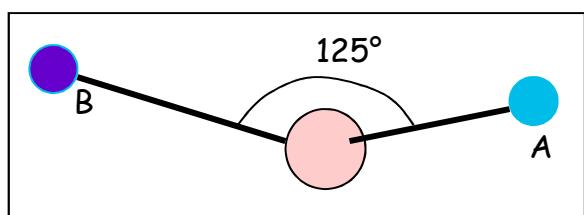
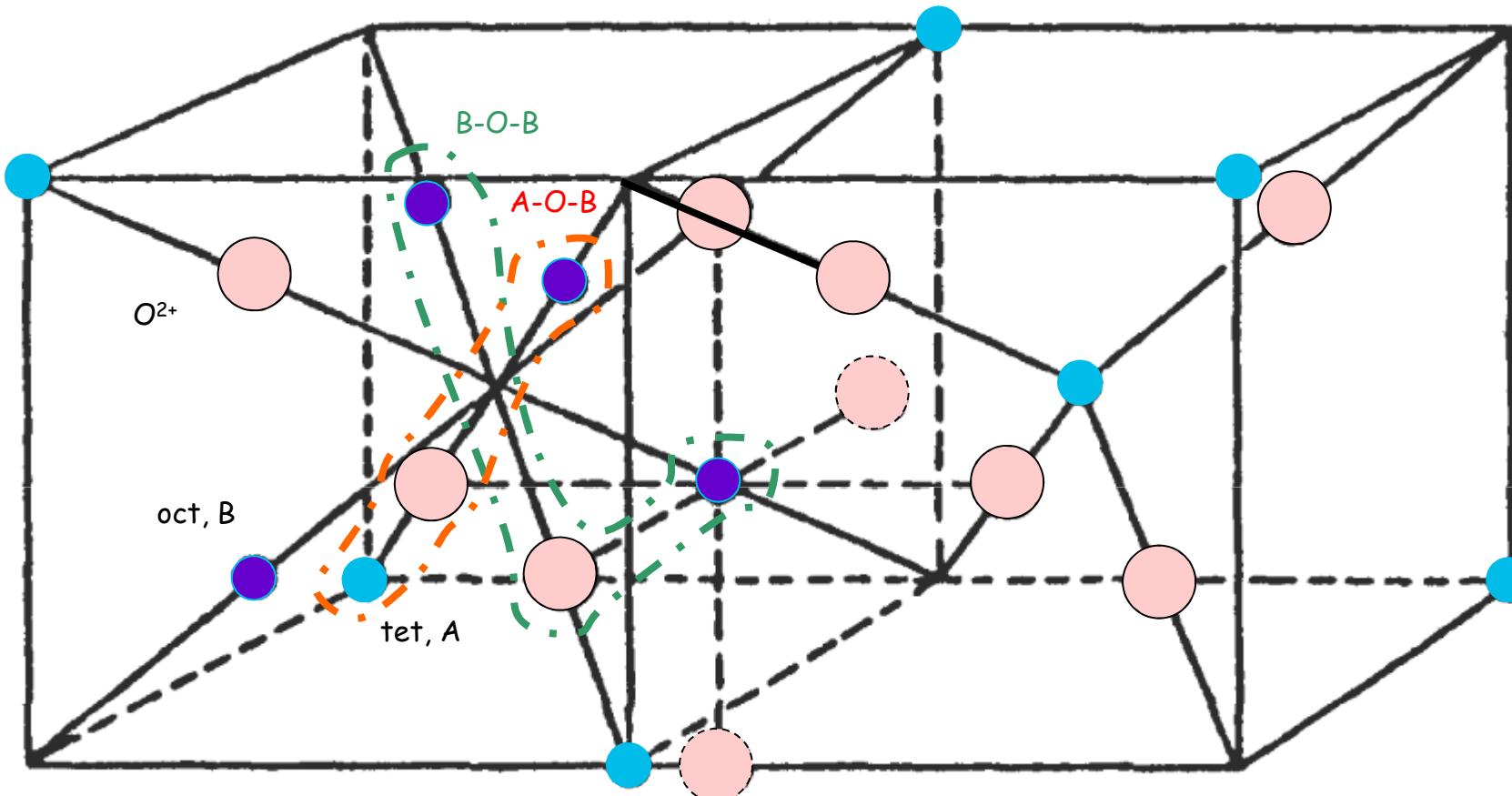
Magnetita



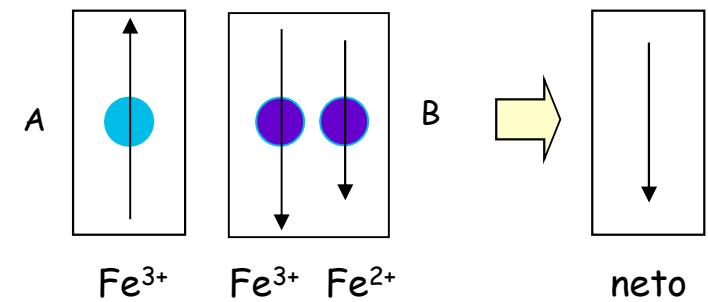
$$\langle \vec{m} \rangle_{Fe_3O_4} \approx m_{Fe^{3+}}(\uparrow\uparrow) + m_{Fe^{3+}}(\downarrow\downarrow) + m_{Fe^{2+}}(\uparrow) \approx m_{Fe^{2+}}(\uparrow) = 4\mu_B \vec{u}_z$$

Intercambio y superintercambio

Magnetita

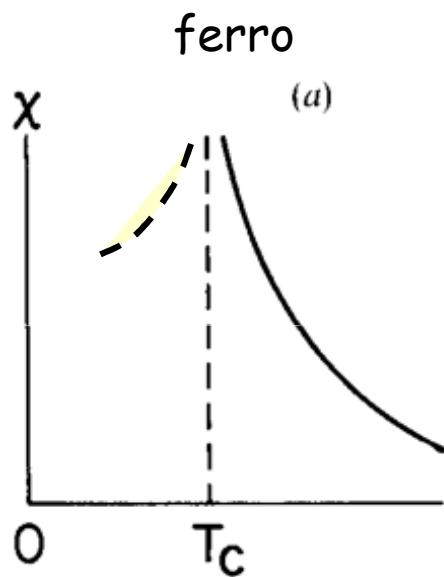


$$|J|_{AB} \gg |J|_{BB}$$



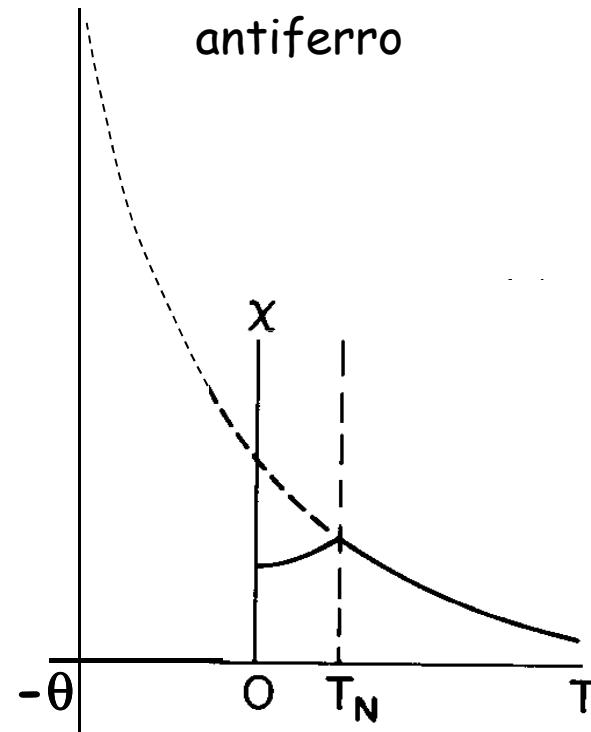
Máximos de χ en T_C , T_N

$$\chi_{inic} = \frac{C}{T + \Theta}$$



$$\chi = \frac{C}{T - T_C}$$

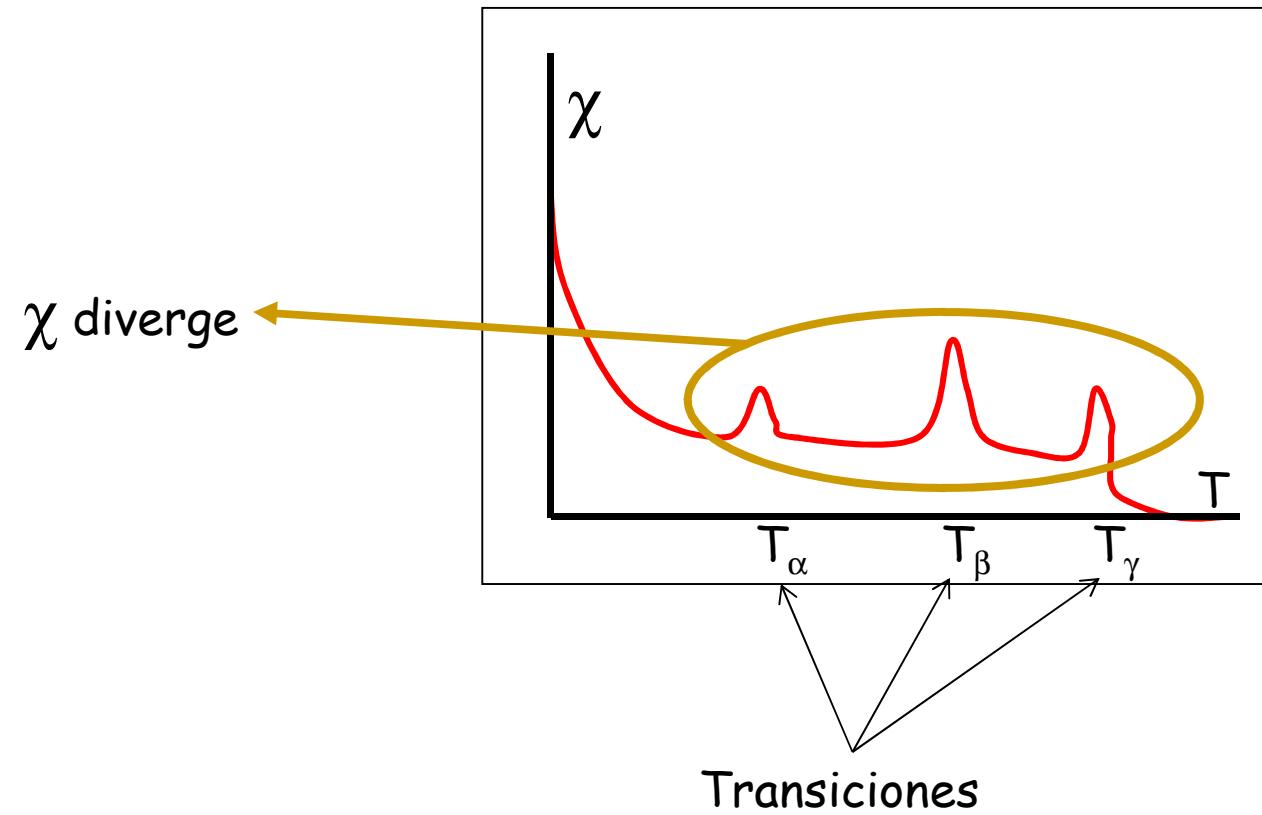
Curie-Weiss Law
($T > T_C$)



$$\chi = \frac{C}{T + \Theta}$$

($T > T_N$)

Comportamiento de χ en puntos de transición de fase magnética



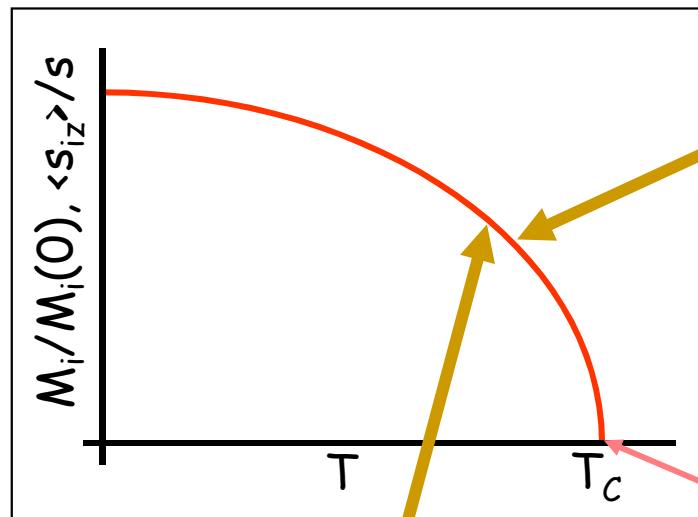
Valores medidos de $m=gs(\mu_B)$ en diferentes materiales

Tabla 3.2. Cristales ferro-ferrimagnéticos

Sustancia	Imanación de saturación		$n_B(0 \text{ K})$	Temperatura de Curie en K
	Temp. amb.	$M_s/10^7 \text{ A/m}$		
Fe	0,1707	0,174	2,22	1043
Co	0,14	0,1446	1,72	1368
Ni	0,0485	0,051	0,606	627
Gd	-	0,206	7,63	292
Dy	-	0,292	10,2	88
MnAs	0,067	0,087	3,4	318
MnBi	0,062	0,068	3,52	630
MnSb	0,071	-	3,5	587
CrO ₂	0,0515	-	2,03	386
MnOFe ₂ O ₃	0,041	-	5,0	573
FeOFe ₂ O ₃	0,048	-	4,1	858
NiOFe ₂ O ₃	0,027	-	2,4	858
CuOFe ₂ O ₃	0,0135	-	1,3	728
MgOFe ₂ O ₃	0,011	-	1,1	713
EuO	-	0,192	6,8	69
Y ₃ Fe ₅ O ₁₂	0,013	0,02	5,0	560

$\langle gs \rangle$

Expresiones empíricas para $\langle s_z \rangle$; Comportamiento crítico



Solución aproximada de la ecuación trascendente - Ferromagnetos

$$\frac{\langle s_z \rangle}{s} = B_s \left(\frac{s}{kT} [\mu_0 g \mu_B H + 2pJ \langle s_z \rangle] \right)$$

$$t = \frac{T}{T_c} \xrightarrow{T < T_c} 1$$

$$\frac{\langle s_z \rangle}{s} = (1-t)^\beta$$

Comportamiento crítico
 $\beta = 3/2 \quad (3d)$

$$\frac{\langle s_z \rangle}{s} = \frac{(1-t)^{\beta'}}{1 - \beta't + At^3 - Ct^{7/2}}$$

Expresión empírica de Aharoni

Introduction to the Theory of Ferromagnetism, Amikam Aharoni, Oxford Science Publications, 1998.

Órdenes de magnitud: Intercambio - Dipolar

Interacción Dipolar

$$J_{dip}, T_{Cr\ dip} ?$$

$$U_{dip} = \frac{\mu_0}{4\pi r^3} [\vec{m}_1 \cdot \vec{m}_2 - 3(\vec{m}_1 \cdot \hat{r})(\vec{m}_2 \cdot \hat{r})]$$

$$U_{dip} \approx \frac{\mu_0}{4\pi r^3} m_1 m_2$$

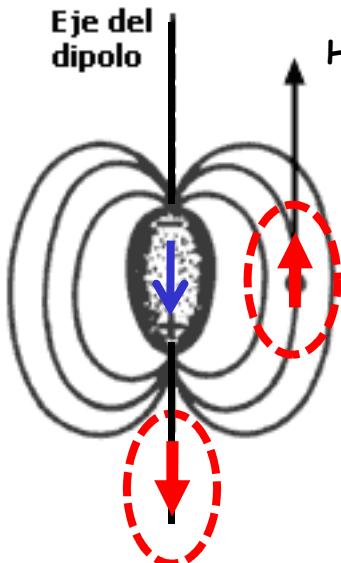
$$m_1 \approx m_2 \approx \mu_B \approx 10^{-23} Am^2$$

$$r \approx 0.2nm = 2 \times 10^{-10} m$$

$$U_{dip} \approx \frac{4\pi \times 10^{-7}}{4\pi \times 8 \times 10^{-30}} \times 10^{-46} Joules = 1.25 \times 10^{-24} Joules = 0.78 \times 10^{-5} eV$$

$$U_{exch} \approx 2J_{exch} = 10^{-21} Joules = 0.63 \times 10^{-2} eV$$

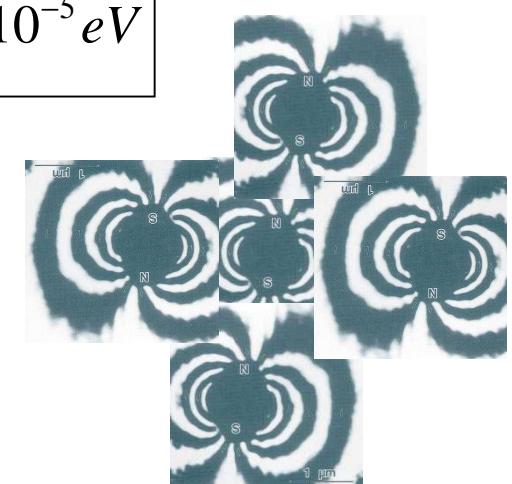
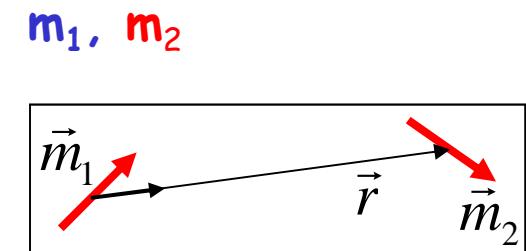
$$\frac{(T_{Cr})_{dip}}{(T_{Cr})_{exch}} \approx \frac{U_{dip}}{U_{exch}} \approx 10^{-3} \xrightarrow{T_{exch} < 10^3 K} (T_{Cr})_{dip} < 1K$$



Intercambio

$$J_{exch} \approx 10^{-21} Joule$$

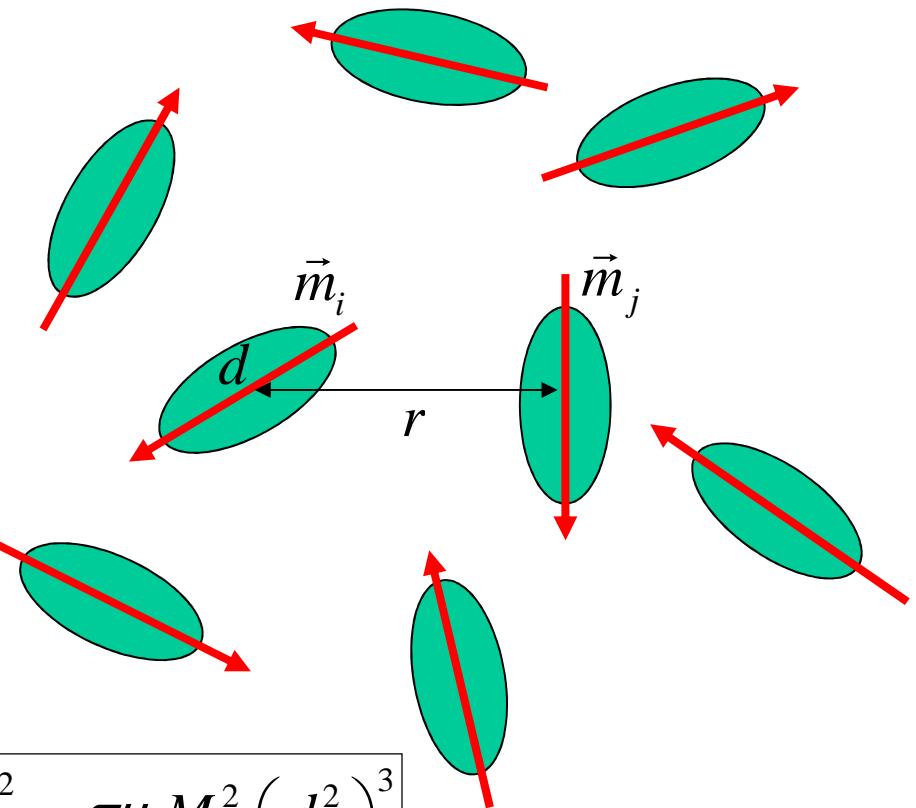
$$T_{Cr\ exch} < 10^3 K$$



Órdenes de magnitud: Dipolar

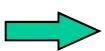
Cómo depende U_{dipolar} de $\langle r \rangle$ y $\langle d \rangle$?

$$|U_{\text{dip}}| \approx \frac{\mu_0}{4\pi r^3} m_i m_j$$



$$|U_{\text{dip}}| \approx \frac{\mu_0 m_i m_j}{4\pi r^3} = \frac{\mu_0 m^2}{4\pi r^3} = \frac{\mu_0 (M_S \pi d^3 / 6)^2}{4\pi r^3} \approx \frac{\pi \mu_0 M_S^2}{144} \left(\frac{d^2}{r} \right)^3$$

Si $d = r$ (partículas en contacto)



$$|U_{\text{dip}}| \approx \frac{\pi \mu_0 M_S^2}{144} d^3$$

Órdenes de magnitud: Dipolar - Zeeman

Para qué relación $\langle r \rangle / \langle d \rangle$ será $U_{\text{Dipolar}} \ll U_{\text{Zeeman}}$?

$$|U_{\text{dip}}| \approx \frac{\mu_0}{4\pi r^3} m_i m_j$$

$$U_z = -\mu_0 \vec{m}_i \cdot \vec{H} \Rightarrow |U_z| \approx \mu_0 m_i H$$

$$\frac{|U_{\text{dip}}|}{|U_z|} \approx \frac{\mu_0 m_i m_j}{4\pi r^3 \mu_0 m_i H} = \frac{m_j}{4\pi r^3 H} = \frac{M_S V_j}{4\pi r^3 H} \approx \frac{M_S 4\pi d^3}{3 \times 4\pi r^3 H} = \frac{M_S}{3H} \left(\frac{d}{r} \right)^3$$

$$\frac{|U_{\text{dip}}|}{|U_z|} \ll 1 \Rightarrow \frac{r}{d} \gg \left(\frac{M_S}{3H} \right)^{1/3} \xrightarrow[M_S=10^6 \text{ Am}^{-1}, H=10^4 \text{ Am}^{-1}]{} \frac{r}{d} \gg 3.2$$

